

The Orchestra of Partial Differential Equations

Adam Larios



19 January 2017

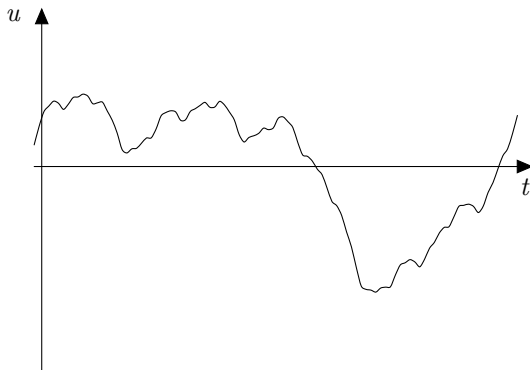
Landscape Seminar

- 1 Fourier Series
- 2 Some Easy Differential Equations
- 3 Some Not-So-Easy Differential Equations

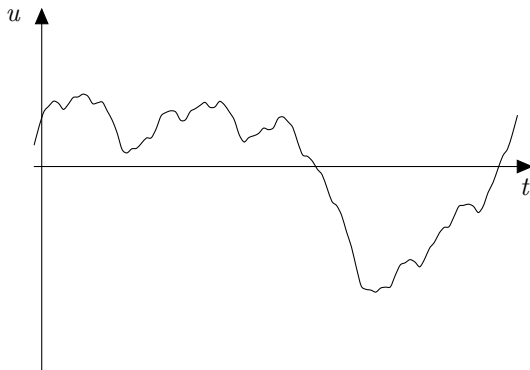
Outline

- 1 Fourier Series
- 2 Some Easy Differential Equations
- 3 Some Not-So-Easy Differential Equations

Frequency

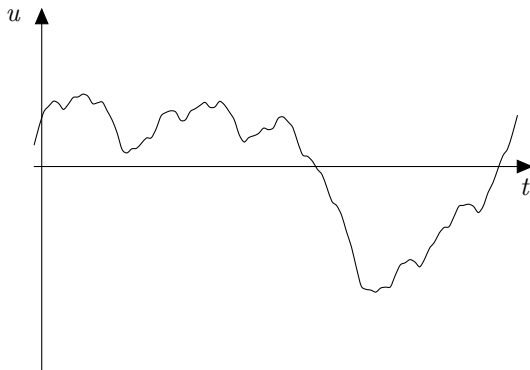


Frequency



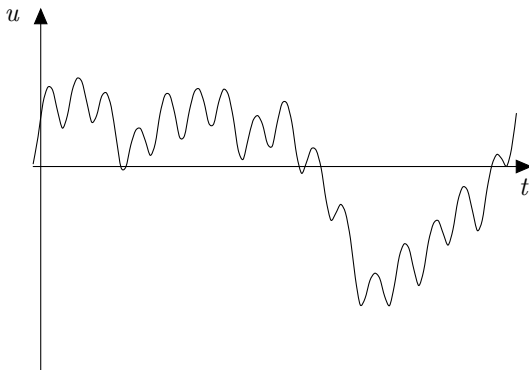
$$u(t) = 0.5 \cos(2t) + 0.125 \cos(8t) + 0.03125 \cos(32t) \\ + 1.0 \sin(1t) + 0.25 \sin(4t) + 0.0625 \sin(16t)$$

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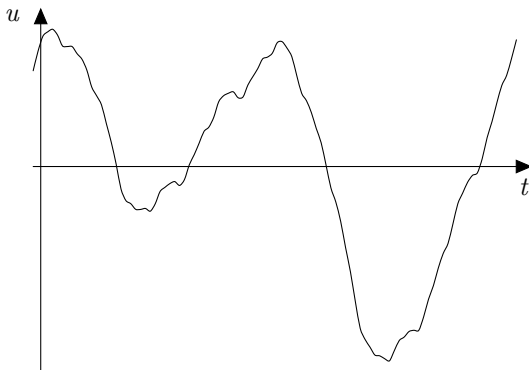
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Fourier Series

$$u(t) = \sum_{k=0}^{\infty} (a_k \cos(kt) + b_k \sin(kt))$$



Fourier Series



$$u(t) = \sum_{k=0}^{\infty} (a_k \cos(kt) + b_k \sin(kt))$$

$$e^{ikt} = \cos(kt) + i \sin(kt)$$

$$\cos(kt) = \frac{e^{ikt} + e^{-ikt}}{2}$$

$$\sin(kt) = \frac{e^{ikt} - e^{-ikt}}{2i}$$

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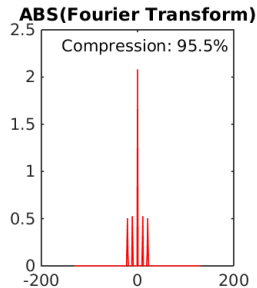
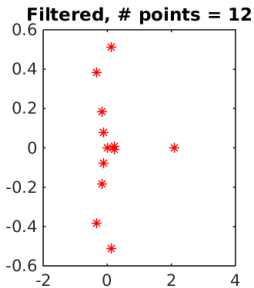
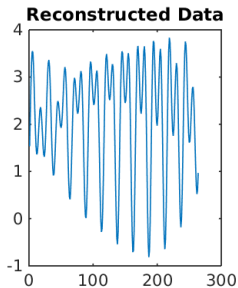
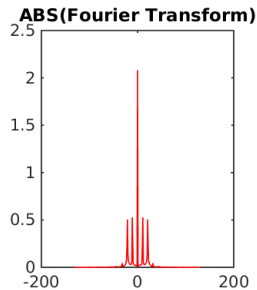
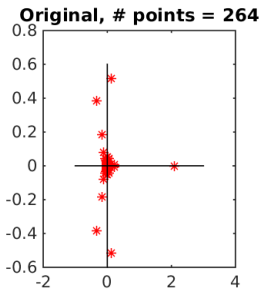
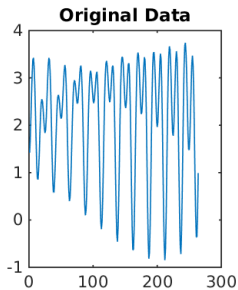
$$\cos(kt) = \frac{e^{ikt} + e^{-ikt}}{2}$$

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$$u(t) = \sum_{k=-\infty}^{\infty} c_k e^{ikt}$$

$$c_k = \frac{1}{2} (a_k - ib_k), \quad k > 0,$$

$$c_k = \frac{1}{2} (a_k + ib_k), \quad k < 0.$$



Multi-dimensional Fourier series

$$u(x) = \sum_{k \in \mathbb{Z}} \widehat{u}_k e^{ikx}$$

$$u(\vec{x}) = \sum_{\vec{k} \in \mathbb{Z}^n} \widehat{u}_{\vec{k}} e^{i\vec{k} \cdot \vec{x}}$$

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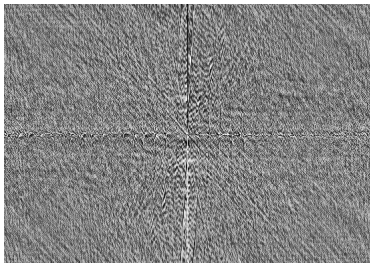
$$\widehat{u}_{\vec{k}} = \frac{1}{2\pi} \int_{-\pi}^{\pi} u(\vec{x}) e^{-i\vec{k} \cdot \vec{x}} d\vec{x}$$



(a) Original image



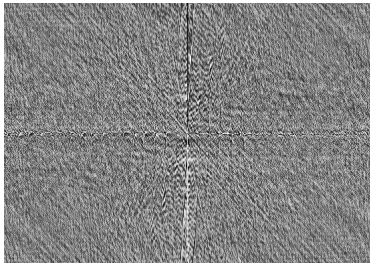
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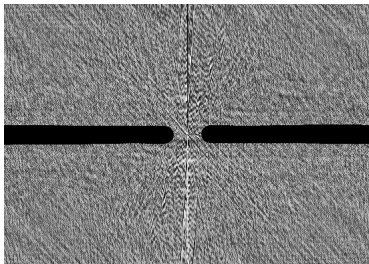
(b) Fourier transform (magnitude)



(a) Original image



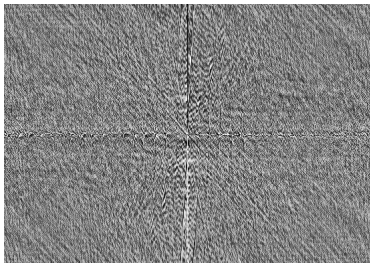
(b) Fourier transform (magnitude)



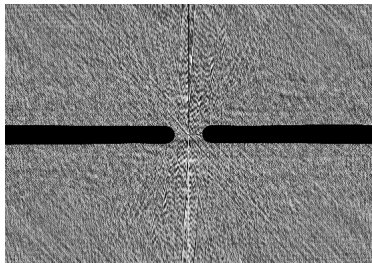
(c) Zero-out Fourier coefficients



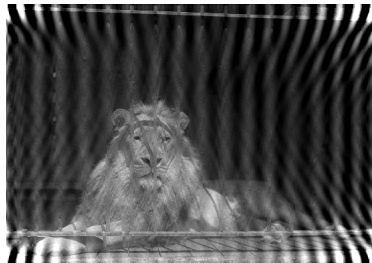
(a) Original image



(b) Fourier transform (magnitude)



(c) Zero-out Fourier coefficients



(d) Inverse Fourier transform

Derivatives

$$u(x) = \sum_{k \in \mathbb{Z}} \hat{u}_k e^{ikx}$$

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Idea

Can the Fourier transform be used to understand differential equations?

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- 2 Some Easy Differential Equations
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The Simplest Partial Differential Equation



$$\frac{d\rho}{dt} =$$

position = $x = x(t)$

velocity = $v = v(t, x(t)) = \frac{dx}{dt}$

density = $\rho = \rho(t, x(t))$

The Simplest Partial Differential Equation



$$\frac{d\rho}{dt} = 0$$

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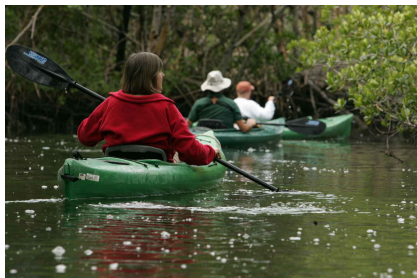
$$\frac{d\rho}{dt} = \frac{\partial\rho}{\partial t} + \frac{dx}{dt} \frac{\partial\rho}{\partial x} = 0$$

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$$\begin{aligned} \frac{d\rho}{dt} &= \frac{\partial\rho}{\partial t} + \frac{dx}{dt} \frac{\partial\rho}{\partial x} \\ &= \frac{\partial\rho}{\partial t} + v \frac{\partial\rho}{\partial x} = 0 \end{aligned}$$

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Transport Equation

$$\rho_t + v\rho_x = 0$$

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Transport Equation

$$\rho_t + v\rho_x = 0$$

Transport Equation in \mathbb{R}^n

$$\rho_t + (\vec{v} \cdot \nabla)\rho = 0$$

$$(\vec{v} \cdot \nabla)\rho = v_1 \frac{\partial\rho}{\partial x} + v_2 \frac{\partial\rho}{\partial y} + v_3 \frac{\partial\rho}{\partial z}$$



Idea

What about the water itself?

What if we set $\rho = v = u$?



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What if we set $\rho = v = u$?

Burgers' Equation

$$u_t + uu_x = 0$$

Burgers' Equation in \mathbb{R}^n

$$\vec{u}_t + (\vec{u} \cdot \nabla) \vec{u} = 0$$

Computer Time!

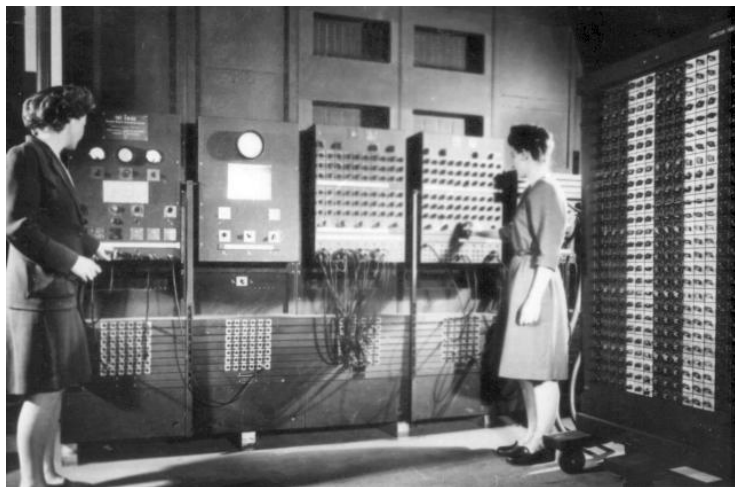


Figure : Programmers working on ENIAC, one of the first computers (c. 1946)

Diffusion Equation



Diffusion

Diffusion Equation



Diffusion

concentration = $\theta = \theta(x, t)$

flux = $f = f(x, t) = f(x)$

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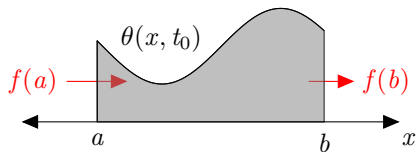
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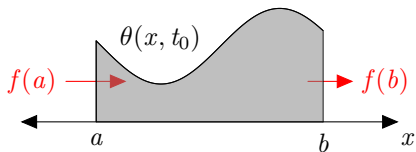
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flux = $f = f(x, t) = f(x)$

$$\frac{d}{dt} \int_a^b \theta(x, t) dx = f(a) - f(b)$$

$$= - \int_a^b \frac{\partial f}{\partial x} dx$$



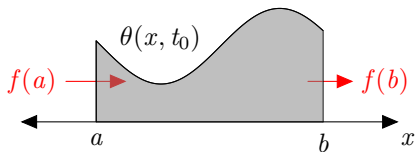
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Fourier's law:

$$f = -\nu \frac{\partial \theta}{\partial x}, \quad \nu > 0$$

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$$\int_a^b \frac{\partial \theta}{\partial t} dx = \int_a^b \nu \frac{\partial^2 \theta}{\partial x^2} dx$$

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Diffusion Equation

$$\theta_t = \nu \theta_{xx}$$

Diffusion Equation in \mathbb{R}^3

$$\theta_t = \nu(\theta_{xx} + \theta_{yy} + \theta_{zz}) = \nu \Delta \theta$$

Diffusion Equation and the Fourier Transform

$$u_t = \nu u_{xx}$$

$$u(x, t) = \sum_{k \in \mathbb{Z}} \hat{u}_k(t) e^{ikx}$$

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Fourier Coefficients

$$\frac{d}{dt} \hat{u}_k = -\nu k^2 \hat{u}_k, \quad k \in \mathbb{Z}$$

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Computer Time Again!



Figure : Woman working on a Cray supercomputer. (c. 1986)

Backward Diffusion

$$u_t = -\nu u_{xx}$$

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$$u(x, t) = \sum_{k \in \mathbb{Z}} e^{+\nu k^2 t} \widehat{u}_k(0) e^{ikx}$$

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Massively unstable!

Fourth-Order Diffusion

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Transport-Diffusion Equation

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$$\rho_t + u\rho_x = \nu\rho_{xx}$$

Transport-Diffusion Equation in \mathbb{R}^n

$$\rho_t + (\vec{u} \cdot \nabla)\rho = \nu\Delta\rho$$

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Nonlinear Equations

Burgers Equation [Shock Waves, Traffic]

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Nonlinear Equations

Burgers Equation [Shock Waves, Traffic]

$$u_t + uu_x = \nu u_{xx}$$

Korteweg-de Vries (KdV) Equation [Water Waves]

$$u_t + uu_x = u_{xxx}$$

Kuramoto-Sivashinsky (KS) Equation [Flames]

$$u_t + uu_x = -\lambda u_{xx} - u_{xxxx}$$

Navier-Stokes Equations [Incompressible Fluids]

$$\begin{cases} \vec{u}_t + (\vec{u} \cdot \nabla) \vec{u} = \nu \Delta \vec{u} - \nabla p \\ \nabla \cdot \vec{u} = 0 \end{cases}$$

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Computer Time Once More!



Figure : Hopper Cray XE6 at NERSC, named after American computer scientist Dr. Grace Hopper, 1906-1992.

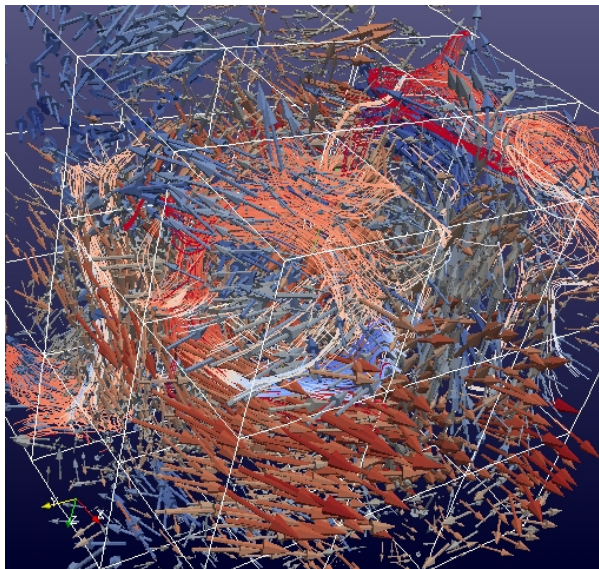
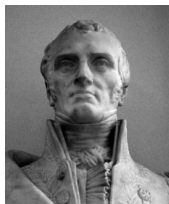


Figure : Simulation of a solution to the 3D Navier-Stokes equations.

The Incompressible Navier-Stokes Equations



Claude L.M.H. Navier

Momentum Equation

$$\underbrace{\frac{\partial \vec{u}}{\partial t}}_{\text{Acceleration}} + \underbrace{(\vec{u} \cdot \nabla) \vec{u}}_{\text{Advection}} = \underbrace{-\nabla p}_{\text{Pressure Gradient}} + \underbrace{\nu \Delta \vec{u}}_{\text{Viscous Diffusion}}$$

Continuity Equation (Divergence-Free Condition)

$$\nabla \cdot \vec{u} = 0$$

Unknowns

$\vec{u} :=$ Velocity (vector)

$p :=$ Pressure (scalar)

Parameter

$\nu :=$ Kinematic Viscosity



George G. Stokes

Problem (Leray 1933)

Existence and uniqueness of strong solutions in 3D for all time. (\$1,000,000 Clay Millennium Prize Problem)

Thank you!